

Mesons in transverse lattice QCD at strong coupling and large- N^* Apoorva D. Patel^a (e-mail: adpatel@cts.iisc.ernet.in)^aCTS and SERC, Indian Institute of Science, Bangalore-560012, India

Mesons in large- N QCD are analysed in light-front coordinates with a transverse lattice at strong coupling. In this limit, their properties can be expressed as simple renormalisations of the 't Hooft model results. The integral eigenvalue equation for the mesons is derived. Spectrum, lightcone wavefunctions and form factors of various mesons can be numerically calculated from this equation.

I have proposed that large- N QCD, with two continuous dimensions on the lightcone and two transverse dimensions on lattice at strong coupling, can be a useful phenomenological model of certain hadronic properties [1]. This combination of large- N and strong coupling limits is chosen to obtain analytically tractable results. The large- N limit allows the use of exact results for two dimensional QCD [2], while the strong coupling limit permits systematic incorporation of the effects of additional transverse directions.

Let μ, ν denote the lightcone directions, and i, j the lattice directions. The action for the theory is (g is held fixed as $N \rightarrow \infty$):

$$\begin{aligned}
S = & \frac{a_{\perp}^2 N}{g^2} \sum_{x_{\perp}} \int d^2 x \left[-\frac{1}{2} \sum_{\mu, \nu} \text{Tr} (F_{\mu\nu}(x) F^{\mu\nu}(x)) \right. \\
& + \bar{\psi}(x) (i \sum_{\mu} \gamma^{\mu} \partial_{\mu} - \sum_{\mu} \gamma^{\mu} A_{\mu} - m) \psi(x) \\
& + i\kappa \sum_i \{ \bar{\psi}(x) (r + \gamma^i) U_i(x) \psi(x + \hat{i} a_{\perp}) \\
& \left. + \bar{\psi}(x) (r - \gamma^i) U_i^{\dagger}(x - \hat{i} a_{\perp}) \psi(x - \hat{i} a_{\perp}) \} \right] \quad (1)
\end{aligned}$$

This action already incorporates $g_{\perp} \rightarrow \infty$, for the gauge coupling corresponding to the transverse directions. The transverse lattice spacing, a_{\perp} , remains finite in this limit. The terms, $\text{Tr}(U_i(x) U_j(x + \hat{i} a_{\perp}) U_i^{\dagger}(x + \hat{j} a_{\perp}) U_j^{\dagger}(x))$ and $\text{Tr}(\partial_{\mu} U_i(x) \partial^{\mu} U_i^{\dagger}(x + \hat{i} a_{\perp}))$, are allowed by symmetry considerations, but they have been dropped because they are suppressed by inverse powers of g_{\perp} . The functional integral over $U_{\perp}(x)$ reduces to the constraint that in any correlation

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function products of $U_{\perp}(x)$ must contract to a colour singlet at each space-time point.

The specification of lattice fermions is not unique, and here I have chosen the nearest neighbour discretisation. $r = 0, 1$ correspond to naive and Wilson fermions respectively. At weak coupling $\kappa = 1/(2a_{\perp})$, but at strong coupling its value is different due to non-perturbative renormalisations. The parameters of the action, that have to be fitted to physical results, are therefore g^2/a_{\perp}^2 , m and κ .

In this theory, the choice of axial gauge $A^+ = 0$ reduces the lightcone gauge field dynamics to a linear confining potential, just as in the case of the 't Hooft model. Unlike the 't Hooft model, however, the γ -matrices cannot be eliminated from fermion propagators and external sources. They become the spin-parity labels of hadron states.

The interplay between large- N and chiral limits is subtle in the 't Hooft model [3]. The chiral symmetry is realised in the Berezinskiĭ-Kosterlitz-Thouless mode in this two-dimensional theory: $\langle \bar{\psi} \psi(x) \bar{\psi} \psi(0) \rangle \sim x^{-1/N}$. The large- N limit must be taken before taking the chiral limit to obtain results relevant to QCD. Contributions from graphs involving non-planar pions are suppressed by inverse powers of N , but are enhanced by inverse powers of m_{π} . If the chiral limit is taken before the large- N limit, one obtains the wrong phase of the theory with massless baryons. If the large- N limit is taken first, then the baryons become infinitely massive solitons and it is not easy to obtain results for them. Properties of mesons can be extracted in a straightforward manner, however, following the analysis of 't Hooft.

The large- N and transverse strong coupling

limits separate the hadron Green functions into two types of non-overlapping factors—propagators along the lightcone and hops along the transverse lattice [1]. The two factors can be evaluated separately and then combined to yield the final results.

The first step is to obtain the renormalised quark propagator. For the quark propagator along the lightcone, the self-energy term consists of two contributions depicted in Fig.1. (a) is the contribution evaluated by 't Hooft,

$$\Sigma_{2d}(p) = -\frac{g^2\gamma^+}{2\pi a_\perp^2 p_-} . \quad (2)$$

(b) is the new contribution due to the transverse lattice—it is a scalar tadpole term, T . Both types of insertions are easily summed up as geometric series. The sum of tadpole insertions renormalises the quark mass, and the sum of gluon corrections shifts the pole position. The quark propagator along the lightcone thus becomes

$$S(p) = \frac{i}{\not{p} - \Sigma_{2d}(p) - m - T} \quad (3)$$

$$= \frac{i(\not{p} - \Sigma_{2d}(p) + m + T)}{2p_+ p_- - (m + T)^2 + (g^2/\pi a_\perp^2)} . \quad (4)$$

The lightcone γ -matrix algebra implies that the part of $S(p)$ proportional to γ^- is the most important for calculations of hadron Green functions.

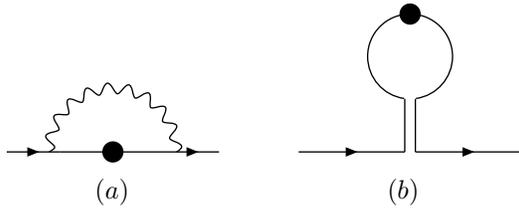


Figure 1. Self energy contributions to the quark propagator: (a) gluonic correction from the continuous dimensions, (b) tadpole correction due to hops in the lattice directions. Fermion line with a filled dot denotes the full quark propagator.

Quark hops along the transverse lattice contribute the factors $i\kappa(r \pm \gamma^i) \exp(\pm ip_i a_\perp)$ to the propagator. These hops are also renormalised by tadpole insertions, which can be taken into account by replacing κ with $\tilde{\kappa} \equiv \kappa/(1 - T_\perp)$. Both

T and T_\perp are calculable functions of the parameters in the action. But they always appear in particular combinations, and so one can parametrise the results in terms of g^2/a_\perp^2 , $\tilde{m} \equiv m + T$ and $\tilde{\kappa}$. (It can be noted that the tadpole corrections vanish for Wilson fermions due to their projection operator spin propagator structure.)

The meson propagator is obtained by including interactions between the renormalised quark and antiquark propagators. There are two types of interactions, as shown in Fig.2. (a) is the gluon exchange in the continuous dimensions, and 't Hooft showed that it produces a linear confining potential between the quark and the antiquark.

$$V_{2d} = -i\frac{g^2}{a_\perp^2} \frac{P}{(k_-)^2} \cdot [\gamma^+ \otimes (\gamma^+)^T] \quad (5)$$

Here k is the gluon momentum and the direct product sign separates the spin factors acting on quark and antiquark propagators. (b) is the new interaction which collapses the quark and the antiquark into a colour singlet meson and then hops the meson along the lattice directions.

$$V_\perp = -\tilde{\kappa}^2 \sum_{i,\pm} \delta(x_\mu - y_\mu) \exp(\pm ip_i a_\perp) \cdot [(r \pm \gamma^i) \otimes (r \mp \gamma^i)^T] , \quad (6)$$

where x and y are quark and antiquark positions, and p is the meson momentum. The position space δ -function constraint translates into a momentum space interaction that is independent of the relative momentum k_μ .

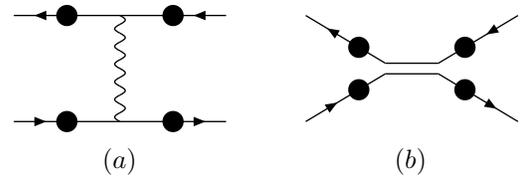


Figure 2. The interactions between quark and anti-quark in a meson: (a) gluon exchange in the continuous dimensions, (b) colour singlet hops in the lattice directions. Fermion line with a filled dot denotes the full quark propagator.

The two types of interactions do not mix with each other, and just produce a series of ladder

diagrams. The properties of the meson poles of the Green functions are easily extracted from a Bethe-Salpeter equation, which sums up the ladder diagrams. For a meson state with spin-parity structure Γ , let

$$\psi_{\Gamma}(p, q) = \langle \bar{\psi}(p - q) \Gamma \psi(q) | \text{Meson}_{\Gamma}(p) \rangle, \quad (7)$$

$$\phi_{\Gamma}(p, q_{-}) = \int d^2 q_{\perp} dq_{+} \psi_{\Gamma}(p, q). \quad (8)$$

Only $\phi_{\Gamma}(p, q_{-})$ appear in the Bethe-Salpeter equation, because the interactions in Eqs.(5,6) do not depend on q_{\perp}, q_{+} . Indeed $\phi_{\Gamma}(p, q_{-})$ are the lightcone wavefunctions of the quarks in the mesons. They satisfy the integral equations

$$\begin{aligned} \Gamma \phi_{\Gamma}(p, q_{-}) &= \int \frac{dq_{+}}{2\pi} S(p - q) \Gamma S(-q) \\ &\cdot \int \frac{dk_{-}}{2\pi} [V_{2d} + V_{\perp}] \phi_{\Gamma}(p, k_{-}). \end{aligned} \quad (9)$$

It is implicitly understood here that the spin factors are ordered along the fermion line. In general, the equations for different spin-parity structures Γ are coupled. The eigenstates can be found by diagonalising the Γ -dependent integral equations, as in the case of conventional strong coupling expansions [4]. Integral equations for mesons with unequal quark masses are obtained by inserting appropriate quark masses in the propagators S and replacing $\tilde{\kappa}^2$ by $\tilde{\kappa}_1 \tilde{\kappa}_2$.

The transverse lattice directions explicitly break the rotational symmetry. Still parity is a good quantum number, and helicity can be defined modulo 4 on the hypercubic lattice. They restrict the mixing amongst various Γ -structures. While perfect rotational symmetry cannot be realised, the free parameters of the theory, g^2/a_{\perp}^2 and $\tilde{\kappa}$, can be non-perturbatively fixed, by demanding that (i) various helicity states belonging to the same angular momentum multiplet be as degenerate as possible, and (ii) the dispersion relation for mesons be as close to the rotationally symmetric situation as possible.

Even without explicitly solving the integral equations, it can be inferred that the meson spectrum consists of a tower of states in each quantum number channel, and the towers of states with different spin form a set of parallel trajectories

[1]. The separations between trajectories depend on the type of transverse lattice discretisation of fermions. Such a behaviour is expected from a successful combination of the 't Hooft model results and strong coupling features.

Explicitly, the integral over q_{+} in Eq.(9) was evaluated by 't Hooft, and the integral of the new interaction V_{\perp} over k_{-} is just a constant. The numerical value of the constant depends on Γ through the commutation properties of various spin factors, and so the meson poles for different Γ -structures are shifted with respect to each other. (From results of conventional strong coupling expansions one can guess that the shifts are essentially independent of the quark masses.)

The actual solutions for the meson spectrum and ϕ_{Γ} have to be obtained numerically. The usefulness of the extreme transverse strong coupling and large- N limit of QCD can only be judged by how well the results fit the experimental data. I expect the formalism to provide a good phenomenological description of deep inelastic scattering, because there the transverse dimensions contribute only at subleading order and so it should not matter much if they are highly distorted by a coarse lattice. The methodology to extract structure functions and form factors from moments of ϕ_{Γ} already exists in the case of 't Hooft model [5,6], and it is straightforward to extend that to the transverse lattice geometry. Detailed numerical analysis of all this is in progress.

REFERENCES

1. A.D. Patel, Nucl. Phys. B (Proc. Suppl.) 94 (2001) 260, hep-lat/0012004.
2. G. 't Hooft, Nucl. Phys. B75 (1974) 461. Here I am following the conventions of: S. Coleman, in "Aspects of Symmetry: Selected Erice Lectures", Cambridge University Press, Cambridge, 1985, p.351.
3. See for example, A.R. Zhitnitsky, Phys. Lett. 165B (1985) 405.
4. R. Friedberg and O. Martin, Nucl. Phys. B279 (1987) 684, and references therein.
5. C.G. Callan, N. Coote and D.J. Gross, Phys. Rev. D13 (1976) 1649.
6. M.B. Einhorn, Phys. Rev. D14 (1976) 3451.