



Diameter rigidity for Kähler manifolds with positive bisectional curvature

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Abstract

We prove that a Kähler manifold with positive bisectional curvature and maximal diameter is isometric to complex projective space with the Fubini-Study metric.

1 Introduction

Let (M, ω) be a Kähler manifold. The *bisectional curvature* of ω along real unit tangent vectors X, Y is defined to be

$$\text{BK}(X, Y) = \text{Rm}(X, JX, JY, Y),$$

where Rm denotes the Riemann curvature tensor of the Riemannian metric associated to ω . In this note we will be concerned with Kähler manifolds (M, ω) satisfying

$$\text{BK} \geq 1, \tag{1}$$

i.e., $\text{BK}(X, Y) \geq 1$ for all real unit tangent vectors X, Y .

A diameter comparison theorem was established for compact Kähler n -manifolds satisfying (1) in [5]. The comparison space here is the complex projective space $\mathbb{C}P^n$

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endowed with the Fubini-Study metric $\omega_{\mathbb{C}P^n}$, normalized so that

$$\int_{\mathbb{C}P^n} \omega_{\mathbb{C}P^n}^n = (2\pi)^n, \text{ equivalently Ric} = (n+1)\omega_{\mathbb{C}P^n}.$$

Theorem 1 (Li-Wang [5]) *If (M^n, ω) is a compact Kähler manifold satisfying $\text{BK} \geq 1$, then*

$$\text{diam}(M) \leq \text{diam}(\mathbb{C}P^n, \omega_{\mathbb{C}P^n}) = \frac{\pi}{\sqrt{2}}.$$

Remark In [5], the diameter bound is stated to be $\pi/2$. This is due to a different normalization for the Hermitian extension of the Riemannian metric.

The main result of this note is a characterization of the case of equality in Theorem 1:

Theorem 2 *Let (M^n, ω) be a compact Kähler manifold satisfying $\text{BK} \geq 1$. If*

$$\text{diam}(M, \omega) = \text{diam}(\mathbb{C}P^n, \omega_{\mathbb{C}P^n}),$$

then (M, ω) is isometric to $(\mathbb{C}P^n, \omega_{\mathbb{C}P^n})$.

The diameter bound in Theorem 1 is analogous to the classical Bonnet-Myers diameter bound for compact Riemannian manifolds with positive Ricci curvature. However, one cannot relax the curvature assumption to a positive Ricci lower bound in the Kähler case, as pointed out in [6]: endow $\mathbb{C}P^1$ with the round metric of curvature $\frac{1}{n+1}$ and consider the product metric on the n -fold product

$$M = \mathbb{C}P^1 \times \dots \times \mathbb{C}P^1.$$

The Ricci curvature of M satisfies $\text{Ric} = (n+1)\omega$, but

$$\text{diam}(M) = \sqrt{\frac{n}{n+1}}\pi > \frac{\pi}{\sqrt{2}},$$

if $n \geq 2$.

In the Riemannian case, the equality case of the Bonnet-Myers diameter bound is the well-known maximal diameter theorem of Cheng. Theorem 2 can be regarded as the Kähler analogue of Cheng's theorem.

Theorem 2 has been established under additional assumptions in [6, 11]. In [6], the authors construct a totally geodesic $\mathbb{C}P^1$ with sectional curvature 2 and use this to show that rigidity holds if $\int_M \omega^n > \pi^n$. In [11], the authors assume that there are compact connected complex submanifolds P and Q in M with $\dim(P) + \dim(Q) = n - 1$ and $d(P, Q) = \frac{\pi}{\sqrt{2}}$. An eigenvalue comparison theorem is then employed to show rigidity.

Our strategy for proving Theorem 2 is to establish a monotonicity formula for a function arising from Lelong numbers of positive currents on $\mathbb{C}P^n$. In [7], the $\partial\bar{\partial}$ -comparison theorem of [11] is reformulated as asserting the positivity of a certain $(1, 1)$ -current and this is the current we work with.

2 Lelong numbers and a monotonicity formula on $\mathbb{C}P^n$

Let M be a Kähler manifold. In what follows, we frequently use the real operator

$$d^c = \frac{\sqrt{-1}}{2\pi}(\bar{\partial} - \partial).$$

Note that

$$dd^c = \frac{1}{\pi}\sqrt{-1}\partial\bar{\partial}.$$

If T is a non-negative current on a M such that

$$T = dd^c\varphi,$$

in a neighbourhood of a point $q \in M$, then the *Lelong number* of T at q is defined as

$$\nu(T, q) := \lim_{r \rightarrow 0^+} \frac{\sup_{B_{\mathbb{C}^n}(0,r)} \varphi(z)}{\log r},$$

where z is a holomorphic coordinate in a neighbourhood of q such that $z(q) = 0$. It is not difficult to see (for instance using the maximum principle) that the quotient on the right is increasing in r , and hence the limit $\nu(T, q)$ exists and is moreover non-negative and independent of the choice of holomorphic coordinates. Note that the normalization is chosen so that if V is a smooth hypersurface with defining function f , and $[V]$ denotes the current of integration along V , then by the Poincaré-Lelong equation, $[V] = dd^c \log |f|$, and so $\nu([V], q) = 1$ for any point $q \in V$.

The following proposition is well known (cf. [2, pg. 164–165]), but since the proof of our main theorem has a precise dependence on the constants involved, we provide a proof for the convenience of the reader.

Proposition 3 *Suppose $T = dd^c\varphi$ as above in a neighbourhood of q with holomorphic coordinates $z = (z^1, \dots, z^n)$ such that $z(q) = 0$. We then have*

$$\nu(T, q) = \lim_{r \rightarrow 0^+} \frac{1}{\pi^{n-1}r^{2n-2}} \int_{B_{\mathbb{C}^n}(0,r)} T \wedge \omega_{\mathbb{C}^n}^{n-1},$$

where $B_{\mathbb{C}^n}(0, r)$ is the ball of radius r around the origin with respect to the Euclidean metric $\omega_{\mathbb{C}^n} = \frac{\sqrt{-1}}{2}\partial\bar{\partial}|z|^2$.

Note that quantity on the right above is increasing in r (cf. [4, pg. 390]), and hence the limit, in particular, exists.

Proof First suppose that φ is smooth. We let

$$v(dd^c \varphi, 0, t) := \frac{1}{\pi^{n-1} t^{2n-2}} \int_{B_{\mathbb{C}^n}(0,t)} dd^c \varphi \wedge \omega_{\mathbb{C}^n}^{n-1},$$

$$\mu_t(\varphi) := \frac{1}{\sigma_{2n-1}} \int_{\mathbb{S}^{2n-1}} \varphi(t, \theta) d\sigma(\theta),$$

where $\sigma_{2n-1} = 2\pi^n / (n - 1)!$ is the volume of the unit sphere in $\mathbb{S}^{2n-1} \subset \mathbb{C}^n$, and $d\sigma$ is the standard Riemannian measure on \mathbb{S}^{2n-1} . Let \mathbb{S}_t^{2n-1} be the sphere of radius t centred at the origin, $d\sigma_t$ the Riemannian measure on it and let $\partial\varphi/\partial\nu$ be the normal derivative of φ . Differentiating in t ,

$$\begin{aligned} \frac{d\mu_t(\varphi)}{dt} &= \frac{1}{\sigma_{2n-1}} \int_{\mathbb{S}^{2n-1}} \frac{\partial\varphi}{\partial t}(t, \theta) d\sigma \\ &= \frac{1}{\sigma_{2n-1} t^{2n-1}} \int_{\mathbb{S}_t^{2n-1}} \frac{\partial\varphi}{\partial\nu} d\sigma_t \\ &= \frac{2}{\sigma_{2n-1} t^{2n-1}} \int_{B_{\mathbb{C}^n}(0,t)} \Delta_{\bar{\partial}} \varphi \frac{\omega_{\mathbb{C}^n}^n}{n!} \\ &= \frac{2}{\sigma_{2n-1} t^{2n-1}} \int_{B_{\mathbb{C}^n}(0,t)} \sqrt{-1} \partial\bar{\partial}\varphi \wedge \frac{\omega_{\mathbb{C}^n}^{n-1}}{(n-1)!} \\ &= \frac{2\pi}{\sigma_{2n-1} (n-1)!} \cdot \frac{1}{t^{2n-1}} \int_{B_{\mathbb{C}^n}(0,t)} dd^c \varphi \wedge \omega_{\mathbb{C}^n}^{n-1} \\ &= \frac{v(T, 0, t)}{t}. \end{aligned}$$

Note that in the third line we have the $\bar{\partial}$ -Laplacian $\Delta_{\bar{\partial}}$, and hence the factor of 2 on application of Green’s formula. Integrating the above equality from r to 1, we obtain the so-called Jensen-Lelong formula (cf. [2, pg. 163]):

$$\mu_1(\varphi) - \mu_r(\varphi) = \int_r^1 v(dd^c \varphi, 0, t) \frac{dt}{t}.$$

By regularization, the above equality also holds for a general, possibly non-smooth, plurisubharmonic function φ . Changing variables $s = \log t$ and dividing by $\log r$ we have

$$\frac{\mu_r(\varphi)}{\log r} = \frac{\mu_1(\varphi)}{\log r} - \frac{1}{\log r} \int_{\log r}^0 v(dd^c \varphi, 0, e^s) ds,$$

and letting $r \rightarrow 0^+$ we obtain

$$\lim_{r \rightarrow 0^+} \nu(T, 0, r) = \lim_{r \rightarrow 0^+} \frac{\mu_r(\varphi)}{\log r}.$$

Next proceeding as in [2, pg. 165], by Harnack inequality and maximum principle, we have that

$$\lim_{r \rightarrow 0^+} \frac{\mu_r(\varphi)}{\log r} = \lim_{r \rightarrow 0^+} \frac{\sup_{z \in \partial B_{\mathbb{C}^n}(0,r)} \varphi(z)}{\log r} = \lim_{r \rightarrow 0^+} \frac{\sup_{z \in B_{\mathbb{C}^n}(0,r)} \varphi(z)}{\log r}.$$

□

We require the following modification, which as far as we can tell, seems to be new.

Proposition 4 *Let T be a non-negative current on $\mathbb{C}P^n$ in a Kähler class, and $q \in \mathbb{C}P^n$. Then*

$$\Theta(T, q, r) := \frac{1}{(2\pi)^{n-1} \sin^{2n-2}(r/\sqrt{2})} \int_{B_{\mathbb{C}P^n}(q,r)} T \wedge \omega_{\mathbb{C}P^n}^{n-1}$$

is increasing in r . Here $B_{\mathbb{C}P^n}(q, r)$ is the ball of radius r with respect to $\omega_{\mathbb{C}P^n}$. Moreover, we also have that

$$\lim_{r \rightarrow 0^+} \Theta(T, q, r) = \nu(T, q). \tag{2}$$

Note that the factor in the denominator is precisely the volume of a ball of radius r in $\mathbb{C}P^{n-1}$ with respect to the Fubini-Study metric $\omega_{\mathbb{C}P^{n-1}}$ upto a factor of $(n - 1)!$.

Proof Let us first assume that T is a smooth $(1, 1)$ Kähler form. We use homogenous coordinates $[\xi_0 : \xi_1 : \dots : \xi_n]$ on $\mathbb{C}P^n$ with $q = [1 : 0 : \dots : 0]$, and the usual in-homogenous coordinates $Z_i = \frac{\xi_i}{\xi_0}$ on $\xi_0 \neq 0$. Then

$$\omega = \sqrt{-1} \partial \bar{\partial} \log |\xi|^2 = \sqrt{-1} \partial \bar{\partial} \log(1 + |Z|^2).$$

We then compute

$$\begin{aligned} \Theta(T, q, r) &= \frac{1}{2^{n-1} \sin^{2n-2}(r/\sqrt{2})} \int_{B_{\mathbb{C}P^n}(q,r)} T \wedge (dd^c \log |\xi|^2)^{n-1} \\ &= \frac{1}{2^{n-1} \sin^{2n-2}(r/\sqrt{2})} \int_{\partial B_{\mathbb{C}P^n}(q,r)} T \wedge d^c \log(1 + |Z|^2) \\ &\quad \wedge (dd^c \log(1 + |Z|^2))^{n-2}. \end{aligned}$$

Now, it is well known fact that

$$\cos^2 \frac{d_{\mathbb{C}P^n}(q, Z)}{\sqrt{2}} = \frac{|\xi_0|^2}{|\xi|^2} = \frac{1}{1 + |Z|^2}.$$

For instance exploiting the $U(n)$ symmetry one needs to check this only for $\mathbb{C}P^1$ which can be done easily. We then have that for any $Z \in \partial B_{\mathbb{C}P^n}(q, r)$,

$$d^c \log(1 + |Z|^2) = \frac{|Z|^2}{1 + |Z|^2} d^c \log |Z|^2 = \sin^2\left(\frac{r}{\sqrt{2}}\right) d^c \log |Z|^2.$$

Putting this back in the formula above we have that

$$\Theta(T, q, r) = \frac{1}{2^{n-1}} \int_{\partial B_{\mathbb{C}P^n}(q,r)} T \wedge d^c \log |Z|^2 \wedge (dd^c \log |Z|^2)^{n-2}. \tag{3}$$

So if $r_1 < r_2$, then integrating by parts we have

$$\Theta(T, q, r_2) - \Theta(T, q, r_1) = \frac{1}{2^{n-1}} \int_{A_{\mathbb{C}P^n}(q,r_1,r_2)} T \wedge (dd^c \log |Z|^2)^{n-1},$$

where $A_{\mathbb{C}P^n}(q, r_1, r_2) = B_{\mathbb{C}P^n}(q, r_2) \setminus \overline{B_{\mathbb{C}P^n}(q, r_1)}$. Now if $\mu : \mathbb{C}P^n \dashrightarrow \mathbb{C}P^{n-1}$ is the projection from q to $[\xi_0 = 0]$, then we have

$$\Theta(T, q, r_2) - \Theta(T, q, r_1) = \frac{1}{(2\pi)^{n-1}} \int_{A_{\mathbb{C}P^n}(q,r_1,r_2)} T \wedge (\mu^* \omega_{\mathbb{C}P^{n-1}})^{n-1} \geq 0.$$

This proves the monotonicity for smooth currents. For a general positive current T we can proceed by regularization. In fact in our case we can first let $r_1 < r_2 < R < \pi/\sqrt{2}$. Then $B(q, R)$ is contained in Euclidean ball (of radius $\tan R$) with respect to the inhomogenous coordinates. We can then use the standard convolution to find sequence of smooth non-negative forms T_j converging weakly to T . Then since $r_1 < r_2 < R$,

$$\Theta(T, q, r_2) - \Theta(T, q, r_1) = \lim_{j \rightarrow \infty} (\Theta(T_j, q, r_2) - \Theta(T_j, q, r_1)) \geq 0.$$

If $r_2 = \pi/\sqrt{2}$, then the result follows by the monotonic convergence.

Next, to compute the limit, we again first work with smooth Kahler forms. If T is smooth then in formula (3), we observe that

$$d^c \log |Z|^2 = \frac{d^c |Z|^2}{|Z|^2} = \frac{d^c |Z|^2}{\tan^2(r/\sqrt{2})},$$

where notice that $d(q, Z) = r$ implies that

$$|Z|^2 = \tan^2\left(\frac{r}{\sqrt{2}}\right).$$

Then we have

$$\Theta(T, q, r) = \frac{1}{2^{n-1}} \int_{\partial B_{\mathbb{C}P^n}(q,r)} T \wedge d^c \log |Z|^2 \wedge (dd^c \log |Z|^2)^{n-2}$$

$$\begin{aligned}
 &= \frac{1}{2^{n-1} \tan^{2n-2}(r/\sqrt{2})} \int_{B_{\mathbb{C}P^n}(q,r)} T \wedge d^c |Z|^2 \wedge (dd^c |Z|^2)^{n-2} \\
 &= \frac{1}{2^{n-1} \tan^{2n-2}(r/\sqrt{2})} \int_{B_{\mathbb{C}P^n}(q,r)} T \wedge (dd^c |Z|^2)^{n-1} \\
 &= \frac{1}{\pi^{n-1} t^{2n-2}} \int_{B_{\mathbb{C}^n}(0,t)} T \wedge \omega_{\mathbb{C}^n}^{n-1},
 \end{aligned}$$

where we integrated by parts in the third line and set $t = \tan(r/\sqrt{2})$, and noted that in terms of the Z -coordinates $B_{\mathbb{C}P^n}(q, r) = B_{\mathbb{C}^n}(0, t)$. Once again by regularization, as above, the above formula holds for general possibly non-smooth currents. Letting $t \rightarrow 0^+$ and applying Proposition 3 we obtain (2). \square

Example 5 (The “model” case) On $\mathbb{C}P^n$ consider the current $T = \sqrt{-1} \partial \bar{\partial} \log |\xi_n|^2 = 2\pi [\xi_n = 0]$, and $q = [1 : 0 : \dots : 0]$. We regard this as the model case for reasons given in Section 3. Then for any $r > 0$,

$$\begin{aligned}
 \int_{B_{\mathbb{C}P^n}(q,r)} T \wedge \omega_{\mathbb{C}P^n}^{n-1} &= 2\pi \int_{B_{\mathbb{C}P^n}(q,r) \cap \{\xi_n=0\}} \omega_{\mathbb{C}P^n}^{n-1} \\
 &= 2\pi \int_{B_{\mathbb{C}P^{n-1}}(q,r)} \omega_{\mathbb{C}P^{n-1}}^{n-1} \\
 &= (2\pi)^n \sin^{2n-2} \left(\frac{r}{\sqrt{2}} \right),
 \end{aligned}$$

and so $\Theta(T, q, r) = 2\pi$ and is independent of r . Note that if we consider a modified

$$\tilde{\Theta}(T, q, r) := \frac{1}{(2\pi)^{n-1} r^{2n-2}} \int_{B_{\mathbb{C}P^n}(q,r)} T \wedge \omega_{\mathbb{C}P^n}^{n-1},$$

where we have r^{2n-2} in the denominator as in the usual Euclidean case, then for T and q as above, we would have that

$$\tilde{\Theta}(T, q, r) = 2\pi \frac{\sin^{2n-2}(r/\sqrt{2})}{r^{2n-2}}.$$

It is easy to see that this function is *decreasing* in r . The increasing property of $\Theta(T, q, r)$ is crucial for our proof of Theorem 2.

3 Proof of the Theorem

In [7], Lott introduces the following current:

$$T_{\omega,p} := \omega + \sqrt{-1} \partial \bar{\partial} \psi_p, \quad \psi_p := \log \cos^2 \left(\frac{d_p}{\sqrt{2}} \right),$$

where p is some fixed point in M and d_p is the distance function from p . Note that *a priori*, $T_{\omega,p}$ is only defined (and also smooth) away from the cut-locus of p . Using the Hessian comparison theorem in [11], which holds away from the cut-locus, Lott observed that T is in fact a global non-negative current if ω satisfies (1).

If $\omega = \omega_{\mathbb{C}P^n}$, and $p = [0 : 0 : \dots : 1]$, then as observed before

$$\cos^2\left(\frac{d_{\omega_{\mathbb{C}P^n},p}}{\sqrt{2}}\right) = \frac{|\xi_n|^2}{|\xi|^2},$$

and so

$$T_{\omega_{\mathbb{C}P^n},p} = \sqrt{-1}\partial\bar{\partial} \log |\xi_n|^2,$$

is precisely the current considered in Example 5 above.

Proof of Theorem First note that by the proof of the Frankel conjecture (cf. [8,10]), M is biholomorphic to $\mathbb{C}P^n$. So from now on we set $M = \mathbb{C}P^n$. Let $p, q \in \mathbb{C}P^n$ such that $d_{\omega,p}(q) = \pi/\sqrt{2}$.

We claim that

$$v(T_{\omega,p}, q) = v(\omega + \pi dd^c \psi_{\omega,p}) \geq 2\pi.$$

To see this, we fix holomorphic coordinates $z := (z^1, \dots, z^n)$ near q with $z(q) = 0$. Then $C^{-1}|z(x)| \leq d(q, x) \leq C|z(x)|$ for some constant $C > 0$, and hence it is enough to show that

$$\lim_{\varepsilon \rightarrow 0^+} \frac{\sup_{B(q,\varepsilon)} \psi_{\omega,p}}{\log \varepsilon} \geq 2,$$

since ω being smooth does not contribute to the Lelong number. It is more convenient to work with

$$\delta_p = \frac{\pi}{2} - \frac{d_p}{\sqrt{2}}.$$

Then $\psi_p = 2 \log \sin \delta_p$. Note that by the diameter upper bound we have $\delta_p(z) \geq 0$ for all z , and that δ_p is Lipschitz with constant $1/\sqrt{2}$. Then for any $x \in \mathbb{C}P^n$,

$$\delta_p(x) \leq \frac{1}{\sqrt{2}}d(q, x),$$

and so $\sup_{B(q,\varepsilon)} \psi_{\omega,p} \leq C + 2 \log \varepsilon$. But then

$$\frac{\sup_{B(q,\varepsilon)} \psi_{\omega,p}}{\log \varepsilon} \geq \frac{C}{\log \varepsilon} + 2 \xrightarrow{\varepsilon \rightarrow 0^+} 2.$$

But then by monotonicity, if $\omega \in c[\omega_{\mathbb{C}P^n}]$, putting $R = \pi/\sqrt{2}$, we have

$$\begin{aligned} 2\pi c &= \frac{1}{(2\pi)^{n-1}} \int_{\mathbb{C}P^n} T \wedge \omega_{\mathbb{C}P^n}^{n-1} = \Theta(T_{\omega,p}, q, R) \geq \lim_{r \rightarrow 0^+} \Theta(T_{\omega,p}, q, r) \\ &= v(T_{\omega,p}, q) \geq 2\pi, \end{aligned}$$

and so $c \geq 1$. On the other hand note that the bisectional curvature lower bound gives

$$\text{Ric}(\omega) \geq (n+1)\omega,$$

and so $c \leq 1$ since $[\text{Ric}(\omega)] = (n+1)[\omega_{\mathbb{C}P^n}]$, and hence $c = 1$. But then the lower bound on the Ricci curvature, and the $\sqrt{-1}\partial\bar{\partial}$ -lemma imply that ω must be Kähler-Einstein and hence isometric to $\omega_{\mathbb{C}P^n}$. \square

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